

eq<sup>n</sup> (4) & (6) become

$$- \frac{dV_s}{dz} = (R + j\omega L) I_s \quad \text{--- (8)}$$

$$- \frac{dI_s}{dz} = (G + j\omega C) V_s \quad \text{--- (9)}$$

Eq<sup>n</sup> (8) & (9) are coupled

Taking second derivative to separate them

$$\frac{d^2 V_s}{dz^2} = (R + j\omega L)(G + j\omega C) V_s$$

$$\Rightarrow \frac{d^2 V_s}{dz^2} - \gamma^2 V_s = 0 \quad \text{--- (10)}$$

where  $\gamma = \alpha + j\beta = \sqrt{(R + j\omega L)(G + j\omega C)}$  --- (11)

taking the second derivative of  $I_s$  in eq<sup>n</sup> (9) and employing eq<sup>n</sup> (8)

$$\frac{d^2 I_s}{dz^2} - \gamma^2 I_s = 0 \quad \text{--- (12)}$$

Eq<sup>n</sup> (10) & (12) are  $\rightarrow$  wave eq<sup>n</sup>s for voltage and current  $\rightarrow$  similar to the wave eq<sup>n</sup>s obtained

for plane waves

$\gamma$  in eq<sup>n</sup> (11)  $\rightarrow$  propagation constant (in per meter)  
 $\alpha \rightarrow$  attenuation constant (in nepers per meter or decibel<sup>2</sup> per meter)  
 $\beta \rightarrow$  phase constant (in radians per meter)

= 6.63 dB



The wavelength  $\lambda$  & wave velocity  $v$  are given by

(31)

$$\lambda = \frac{2\pi}{\beta} \quad , \quad \boxed{v = \frac{\omega}{\beta} = f\lambda} \quad \text{---(14)}$$

Sol<sup>n</sup>s of homogeneous differential eq<sup>n</sup> (10) & (12)

$$V_s(z) = V_0^+ e^{-\gamma z} + V_0^- e^{+\gamma z} \quad \text{---(15)}$$

and

$$I_s(z) = I_0^+ e^{-\gamma z} + I_0^- e^{+\gamma z} \quad \text{---(16)}$$

$V_0^+, V_0^-, I_0^+, I_0^-$   $\rightarrow$  wave amplitude  
 $+, -$   $\rightarrow$  wave traveling along  $+z$  and  $-z$  directions.

$\Rightarrow$  Instantaneous expressions for voltage as

$$\begin{aligned} V(z,t) &= \text{Re} [V_s(z) e^{j\omega t}] \\ &= V_0^+ e^{-\alpha z} \cos(\omega t - \beta z) + V_0^- e^{+\alpha z} \cos(\omega t + \beta z) \end{aligned} \quad \text{---(17)}$$

Characteristic impedance  $Z_0$  of the line

$\hookrightarrow$  ratio of positively traveling voltage wave to the current wave at any point on the line.

$Z_0 \rightarrow$  analogous to the  $\eta \rightarrow$  intrinsic impedance of the medium of wave propagation.

Substituting eq<sup>n</sup> (15) and (16) in eq<sup>n</sup> (8) & (9) and equating coefficients of  $e^{+\gamma z}$  and  $e^{-\gamma z}$

$$Z_0 = \frac{V_0^+}{I_0^+} = -\frac{V_0^-}{I_0^-} = \frac{R + j\omega L}{\gamma} = \frac{\gamma}{\gamma + j\omega C} \quad \dots (12)$$

or

$$Z_0 = \frac{\sqrt{R + j\omega L}}{\sqrt{\gamma + j\omega C}} = R_0 + j X_0 \quad \dots (13)$$

Real part  $\downarrow$  ohmsmeter  
Imaginary part  $\downarrow$  ohmsmeter

$R_0$   
 $\downarrow$   
ohms

$\gamma$   
 $Z_0$  }  $\rightarrow$  important properties of the line

$\hookrightarrow$  Both depend on the line parameters  $R, L, C_1$  and  $C$  and frequency of operation.

admittance  $Y_0 = \frac{1}{Z_0}$ .

We have considered lossy type transmission line  $\rightarrow$  The conductors comprising the line are imperfect ( $\sigma_c \neq \infty$ ) and the dielectric in which the conductors are embedded is lossy ( $\sigma \neq 0$ )

# Lossless Line ( $R=0=G$ )

A transmission line  $\rightarrow$  lossless  
 if the conductors of the line are perfect ( $\sigma_c \approx \infty$ )  
 and the dielectric medium separating them is  
 lossless ( $\sigma = 0$ )  
 when  $\sigma_c \approx \infty, \sigma \approx 0$

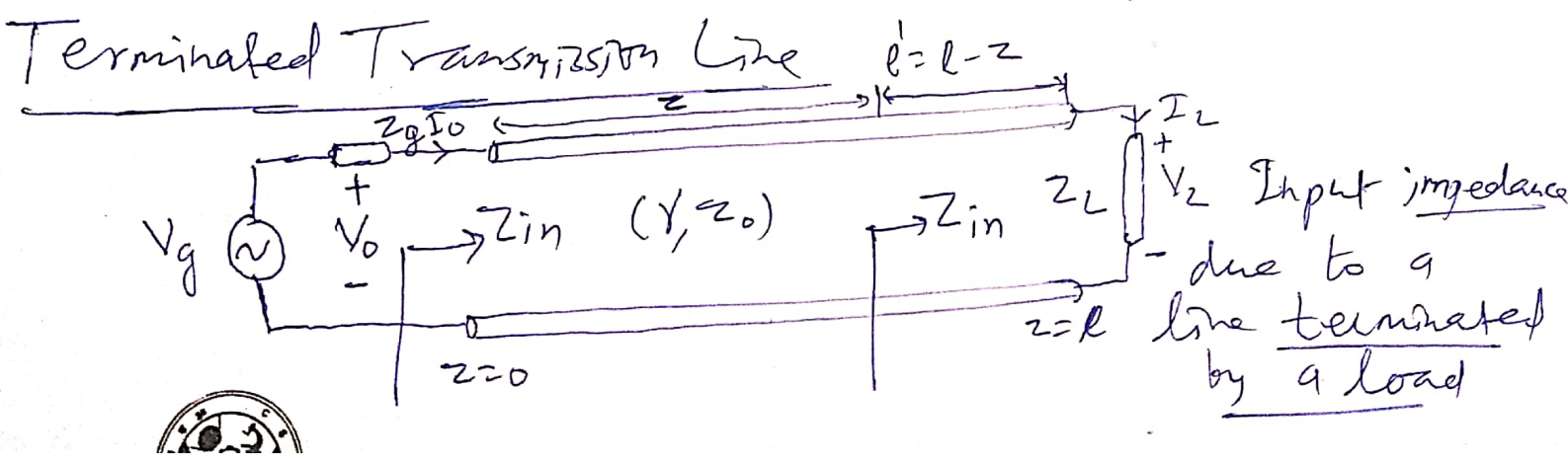
$$\boxed{R=0=G} \quad \text{--- (20)}$$

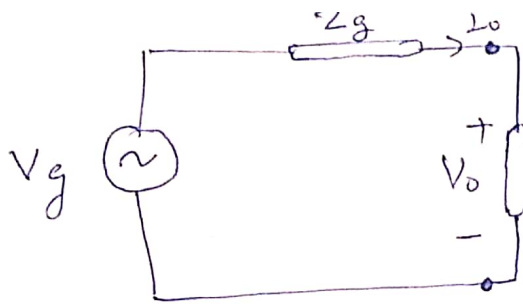
$\downarrow$   
 Necessary conditions for a line to be lossless.  
 For such a line, eq<sup>n</sup> (20) forces eq<sup>s</sup> (11), (14) and (15)  
 to become

$$\alpha = 0, \quad \gamma = j\beta = j\omega\sqrt{LC} \quad \text{--- 21 (a)}$$

$$v = \frac{\omega}{\beta} = \frac{1}{\sqrt{LC}} = f\lambda \quad \text{--- 21 (b)}$$

$$X_0 = 0, \quad Z_0 = R_0 = \sqrt{\frac{L}{C}} \quad \text{--- 21 (c)}$$





(34)

Equivalent circuit for finding  $V_0$  and  $I_0$  in terms of  $Z_{in}$  at input

Transmission line of length  $l$  characterised by  $\gamma$  and  $Z_0$ , connected to a load  $Z_L$ . The generator sees the line with the load as an input impedance  $Z_{in}$ .

Transmission line extend from  $z=0$  at the generator to  $z=l$  at the load.

The voltage and current eq<sup>n</sup> in (15) & (16)

incorporating eq<sup>n</sup> (18)

$$V_s(z) = V_0^+ e^{-\gamma z} + V_0^- e^{\gamma z} \quad \text{--- (22)}$$

$$I_s(z) = \frac{V_0^+}{Z_0} e^{-\gamma z} - \frac{V_0^-}{Z_0} e^{\gamma z} \quad \text{--- (23)}$$

$V_0^+$   
 $V_0^-$

} Terminal conditions must be given

if we are given the conditions at the input

say  $V_0 = V(z=0), \quad I_0 = I(z=0) \quad \text{--- (24)}$

Substituting it into eq<sup>n</sup> (22) (23)

$$V_0^+ = \frac{1}{2} (V_0 + Z_0 I_0) \quad \text{--- 25 (a)}$$

$$V_0^- = \frac{1}{2} (V_0 - Z_0 I_0) \quad \text{--- 25 (b)}$$

If the input impedance at the input terminal is  $Z_{in}$ , the input voltage  $V_0$  and the input current  $I_0$  are obtained as

$$V_0 = \frac{Z_{in}}{Z_{in} + Z_g} V_g, \quad I_0 = \frac{V_g}{Z_{in} + Z_g} \quad (26)$$

If we are given the conditions at the load, say

$$V_L = V(z=l), \quad I_L = I(z=l) \quad (27)$$

from eq<sup>n</sup> (22) & (23)

$$V_0^+ = \frac{1}{2} (V_L + Z_0 I_L) e^{\gamma l} \quad (28a)$$

$$V_0^- = \frac{1}{2} (V_L - Z_0 I_L) e^{-\gamma l} \quad (28b)$$

We find the input impedance  $Z_{in} = V_s(z)/I_s(z)$  at any point on the line. At the generator, for example eq<sup>n</sup> (28) and (23) yield

$$Z_{in} = \frac{V_s(z)}{I_s(z)} = \frac{Z_0 (V_0^+ + V_0^-)}{V_0^+ + V_0^-} \quad (29)$$

from eq<sup>n</sup> (28) & (29) and

$$\frac{e^{\gamma l} + e^{-\gamma l}}{2} = \cosh \gamma l; \quad \frac{e^{\gamma l} - e^{-\gamma l}}{2} = \sinh \gamma l \quad (30a)$$

$$\text{or} \quad \tanh \gamma l = \frac{\sinh \gamma l}{\cosh \gamma l} = \frac{e^{\gamma l} - e^{-\gamma l}}{e^{\gamma l} + e^{-\gamma l}} \quad (30b)$$

$$\Rightarrow Z_{in} = Z_0 \left[ \frac{Z_L + Z_0 \tanh \gamma l}{Z_0 + Z_L \tanh \gamma l} \right] \text{ (lossy)} \quad (31)$$